

Chapter 11Dual Nature of Radiation & Matter

Light behaves both as a wave (interference, diffraction) and as a stream of particles called ~~quanta~~ photons ($E = h \nu$).

Big question :

If light has particle nature, can matter* (electron, proton...) have wave nature?

de Broglie's bold answer : YES.

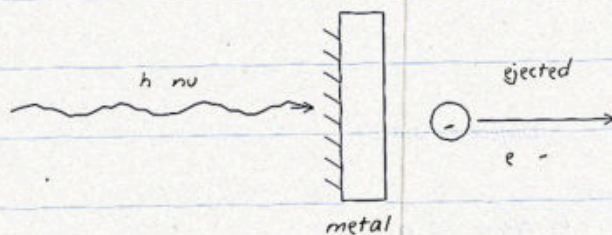


Fig. photon strikes metal \rightarrow electron is ejected

Topics ahead

1. Electron emission ; work function
2. Photoelectric effect ; Einstein eqn.
3. Matter waves ; Davisson - Germer.

Electron Emission

Metals contain free electrons . But at room temperature they cannot escape the surface .
A potential barrier holds them in.

Work function W

Minimum energy needed by a free electron to just escape the metal surface .

$$W = h \nu_0$$

<- threshold
<- frequency

*

Unit : joule (J) or electron - volt (eV)

$$1 \text{ eV} = 1.6 \times 10^{-19} \text{ J}$$

Typical W values

Caesium (Cs) : 2.14 eV

Potassium : 2.30 eV

Sodium : 2.75 eV

Zinc : 3.60 eV

Copper : 4.65 eV

Platinum : 5.65 eV

Smaller W \rightarrow easier to eject electrons ;

Cs , K , Na are used in photo - cells.

(low W matches visible light frequencies.)

Types of Electron Emission

Energy to overcome W can be supplied in four physical ways :

(1) Thermionic emission

Heat the metal . Free electrons gain thermal energy and escape .

Used in : CRT cathode , vacuum ~~diode~~ valves.

(2) Photoelectric emission

Light of suitable frequency falls on metal .

Photons transfer energy to electrons .

Used in : photocells , solar cells.

(3) Field (cold) emission *

A very strong external electric field

(10^8 V/m) pulls electrons out of the surface (no heating needed).

(4) Secondary emission

Fast incoming particles knock out further electrons from the metal surface .

Used in : photomultiplier tubes.

All four overcome the same barrier = W .

Photoelectric Effect

Emission of electrons from a metal surface when light of suitable frequency falls on it.

Historical background

Hertz (1887) : UV light on Zn knob made sparks jump more easily across a gap.

Hallwachs & Lenard (1886 - 1902) studied the effect systematically :

- * Negative Zn plate lost charge in UV light.
- * Positive Zn plate did NOT lose charge.
- * So UV ejected negative particles (electrons).

Key features (observed)

1. Below a threshold frequency ν_0 , no emission - whatever the intensity.
2. Above ν_0 , emission is instantaneous ($< 10^{-9}$ s).
3. Photocurrent is proportional to intensity.
4. Max KE of photoelectrons depends only on ν , NOT on intensity.

Classical wave theory cannot explain these.

Experimental Setup

Evacuated glass tube with a quartz window
(UV passes through quartz, not ordinary glass).

Two electrodes : emitter C , collector A.

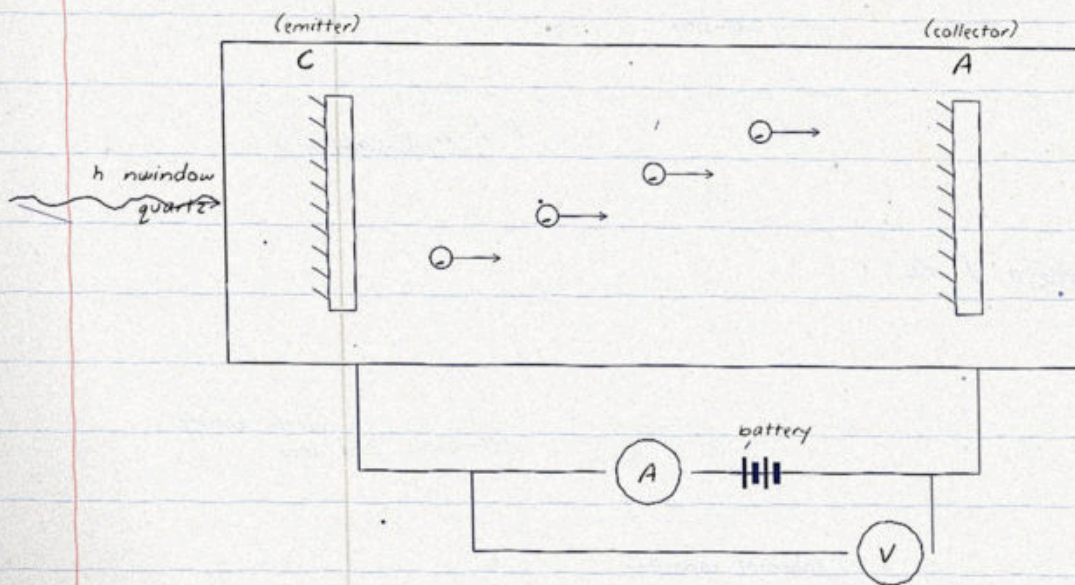


Fig. photoelectric experiment apparatus

Ammeter A reads photocurrent i .

Voltmeter V reads p.d. between C and A.

Polarity can be reversed by a commutator.

Effect of Intensity on Current

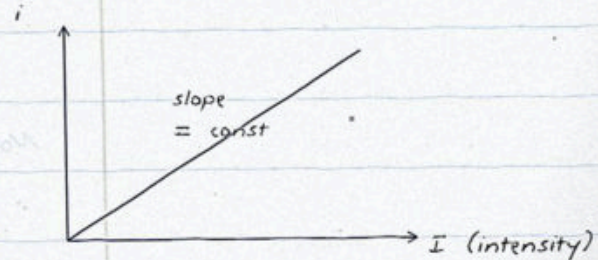
Keep frequency ν and collector p.d. fixed .

Vary intensity I of incident light.

Observation :

Photocurrent i is directly proportional to the intensity of incident light.

i proportional to I ← linear



Why ? (photon view)

Higher intensity \Rightarrow more photons per second.

More photons \rightarrow more electrons ejected per s.

But each photon still carries energy $E = h \nu$. . .

Hence intensity controls NUMBER, not energy of emitted photoelectrons.

Effect of Potential : $V - I$ Curve

Keep ν and I (intensity) fixed. Vary the potential of collector A w.r.t. emitter C.

Source: *Handwritten*

*

Positive V : i rises, reaches saturation.

Negative V : i falls; at $V = -V_0$ $i = 0$.

V_0 is called the stopping potential.

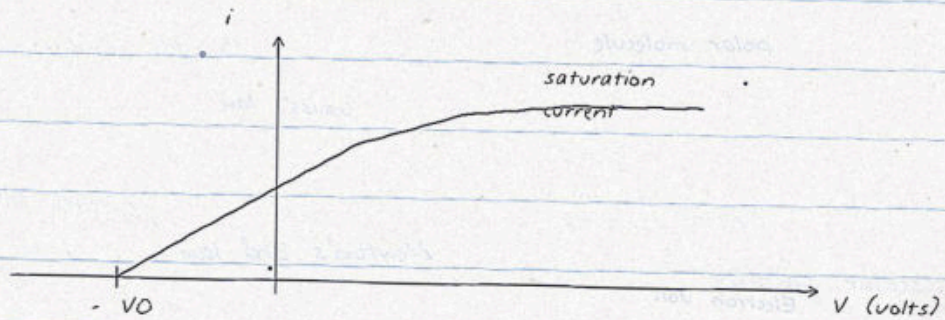


Fig. $V - I$ characteristic of photocell

Saturation $i \rightarrow$ all emitted electrons reach A.

At $V = -V_0$ even the fastest electron is stopped just before reaching A.

Stopping Potential V_0

Definition : the negative collector voltage that just stops the fastest photoelectron .

— Energy balance :

$$e V_0 = (1/2) m v_{\max}^2 = KE_{\max} \quad \begin{array}{l} \leftarrow \text{definition} \\ \leftarrow \text{of } V_0 \end{array}$$

Three key results

① V_0 is independent of intensity.

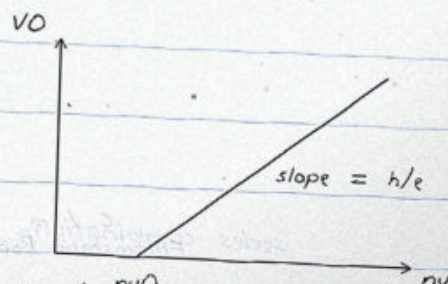
More photons, but same KE per electron.

② V_0 increases linearly with ν .

$$V_0 = (h/e) \nu - W/e$$

③ V_0 is zero at $\nu = \nu_0$.

Below ν_0 no emission at all .



Slope of V_0 vs. ν line = h/e .

Failure of Wave Theory

Classical wave theory predicts that energy of a wave is shared by the whole wavefront and depends only on intensity (amplitude 2).

Contradiction (1)

Wave view : enough intensity should always eject electrons, regardless of frequency.

* Actual obs. : no emission below ~~intensity~~ ν_0 .

Contradiction (2)

Wave view : weak light should take time to build up energy; emission delayed by seconds. *

Actual obs. : emission is instantaneous (10^{-9} s).

Contradiction (3)

Wave view : KE of electron should grow with intensity.

Actual obs. : KE_{max} depends only on ν .

Conclusion

Light energy comes in discrete packets, each interacting with ONE electron at a time.

Einstein (1905) : introduce photons.

Einstein's Photoelectric Equation

Light of frequency ν is made of photons, each of energy $E = h \nu$. One photon transfers ALL its energy to a single electron.

Energy balance on the surface :

$$h \nu = W + \frac{1}{2} m v_{\max}^2$$

$$h \nu = W + KE_{\max}$$

<- Einstein
<- 1905

Substituting $KE_{\max} = e V_0$:

$$h \nu = W + e V_0$$

Special case : $\nu = \nu_0$

Then $KE_{\max} = 0$, so :

$$W = h \nu_0$$

$$\text{i.e. } \nu_0 = W / h$$

Below ν_0 , KE_{\max} would be negative - impossible, so no emission occurs.

This is the meaning of threshold frequency.

Stopping-Potential Formula

From $h\nu = W + eV_0$:

$$V_0 = (h/e) \cdot \nu - W/e$$

<- linear in ν

Compare with $y = m x + c$:

$$\text{slope } m = h/e$$

$$\text{intercept } c = -W/e$$

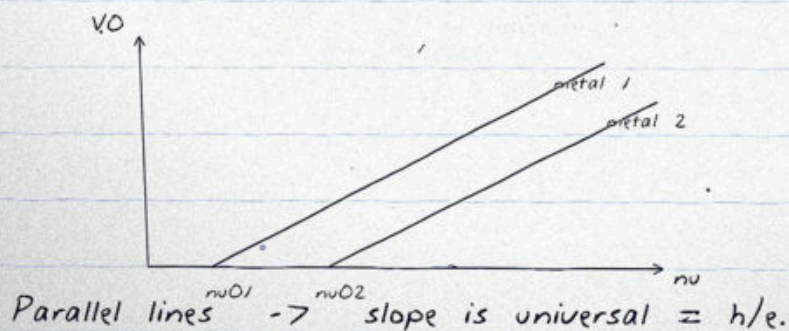
$$x\text{-intercept} = \nu_0 = W/h$$

Millikan's Verification (1916)

Plotted V_0 vs ν for several metals .

Lines were parallel , slope $= h/e$.

Gave $h = 6.57 \times 10^{-34} \text{ J s}$ (good match).



Photon - the Particle of Light

Photon = quantum of EM radiation .

Each photon carries definite energy and momentum^{*} , travels at speed of light c .

Energy of a photon

$$E = h \nu = h c / \lambda$$

ν ← frequency
 λ ← or wavelength

Momentum of a photon

Photon is massless but carries momentum :

$$p = E / c = h \nu / c = h / \lambda$$

Properties of photon

- * Rest mass = 0 (only relativistic mass).
- * Electric charge = 0 , unaffected by E , B .
- * Energy of photon does not change with intensity - only the number of photons.
- * Photons follow ~~Newton~~ Bose - Einstein statistics.
- * Interaction : all - or - nothing with one e^- .

Photon has wave character (ν , λ) too.

Worked Example 1

Q. Light of wavelength 488 nm hits a Cs plate ($W = 2.14$ eV). Find KE_{\max} of the photoelectron and stopping potential V_0 .

Sol :

Photon energy $E = hc / \lambda$

$$E = (6.63 \times 10^{-34}) (3 \times 10^8) / (488 \times 10^{-9})$$

$$= 4.08 \times 10^{-19} \text{ J}$$

$$= 4.08 / 1.6 \text{ eV} = 2.55 \text{ eV}$$

$$KE_{\max} = h\nu - W = E - W \\ = 2.55 - 2.14 = 0.41 \text{ eV}$$

$$KE_{\max} = 0.41 \text{ eV}$$

Stopping potential $V_0 = KE_{\max} / e$

$$V_0 = 0.41 \text{ eV} / e = 0.41 \text{ V}$$

$$V_0 = 0.41 \text{ V}$$

Worked Example 2

Q. Two photoelectric measurements on a metal :

$$\nu_1 = 6.0 \times 10^{14} \text{ Hz}, \quad V_{0_1} = 0.42 \text{ V}$$

$$\nu_2 = 7.5 \times 10^{14} \text{ Hz}, \quad V_{0_2} = 1.04 \text{ V}$$

Find h , then W and ν_0 .

Sol :

$$h \nu = W + e V_0 \quad \text{for both readings.}$$

$$\text{Subtract : } h(\nu_2 - \nu_1) = e(V_{0_2} - V_{0_1})$$

$$h = e(V_{0_2} - V_{0_1}) / (\nu_2 - \nu_1)$$

$$= (1.6 \times 10^{-19})(1.04 - 0.42)$$

$$/ (7.5 - 6.0) \times 10^{14}$$

$$= (1.6 \times 0.62) / 1.5 \times 10^{-33}$$

$$\text{approx } 6.6 \times 10^{-34} \text{ J s. } (= h)$$

$$\text{Now } W = h \nu_1 - e V_{0_1}$$

$$= (6.6 \times 10^{-34})(6.0 \times 10^{14})$$

$$- (1.6 \times 10^{-19})(0.42)$$

$$= 3.96 \times 10^{-19} - 0.67 \times 10^{-19}$$

$$\text{approx } 3.29 \times 10^{-19} \text{ J} = 2.06 \text{ eV.}$$

$$W = 2.06 \text{ eV}, \quad \nu_0 = W/h = 5.0 \times 10^{14} \text{ Hz}$$

Visible light (red - violet) can eject electrons.

Wave Nature of Matter

Photons (mass = 0) have momentum $p = h / \lambda$.

L. de Broglie (1924) : why not extend this to all material particles ?

de Broglie hypothesis

A moving particle of momentum p has a wave associated with it, of wavelength :

$$\lambda = h / p = h / (m v) \quad \begin{array}{l} \leftarrow \text{matter} \\ \leftarrow \text{wave} \end{array}$$

Implication : λ is small for heavy objects (large p), large for light particles.

From energy

IF $KE = (1/2) m v^2$, then $p = \sqrt{2 m KE}$.

$$\lambda_* = h / \sqrt{2 m KE} \quad *$$

* For a charged particle accelerated through potential V : $KE = q V$.

$$\lambda = h / \sqrt{2 m q V}$$

de Broglie Wave Picture

Matter wave is a probability wave, not a mechanical wave. Amplitude squared \rightarrow probability of finding the particle.

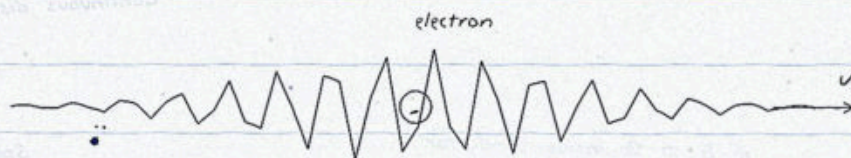


Fig. matter wave packet associated with a moving e^-

Macro vs Micro

- * Heavy objects (m large) \rightarrow λ tiny, wave behaviour undetectable.
- * Light particles (e^- , p , n) \rightarrow λ atomic spacing \rightarrow diffraction observed.

λ for charged particles depends on charge and mass, not on the medium.

All particles have wave nature (de Broglie).

Electron Wavelength Formula

An electron is accelerated from rest through a potential difference V .

$$\text{Energy gained} = (1/2) m v^2 = e V$$

$$v = \sqrt{2 e V / m}$$

$$p = m v = \sqrt{2 m e V}$$

$$\lambda = h / \sqrt{2 m e V} \quad \leftarrow \text{exact}$$

Putting numbers (h, m, e in SI) :

$$\lambda = (6.63 \times 10^{-34}) /$$

$$\sqrt{2 \times 9.1 \times 10^{-31} \times 1.6 \times 10^{-19} V}$$

$$= 1.227 \times 10^{-9} / \sqrt{V} \text{ m}$$

$$= 12.27 / \sqrt{V} \text{ angstrom}$$

$$\lambda \text{ (in \AA)} = 12.27 / \sqrt{V}$$

\leftarrow useful
 \leftarrow shortcut

Sample values

$$V = 100 \text{ V} \quad \rightarrow \quad \lambda = 1.23 \text{ \AA}$$

$$V = 1000 \text{ V} \quad \rightarrow \quad \lambda = 0.39 \text{ \AA}$$

$$V = 10 \text{ kV} \quad \rightarrow \quad \lambda = 0.12 \text{ \AA}$$

(comparable to atomic spacings $\sim 1 \text{ \AA}$).

Worked Example 3

Q. Find de Broglie wavelength of :

(a) electron with $v = 5.4 \times 10^6 \text{ m/s}$

(b) cricket ball $m = 0.15 \text{ kg}$, $v = 20 \text{ m/s}$

(c) proton $v = 1.0 \times 10^5 \text{ m/s}$

Sol :

(a) $m_e = 9.11 \times 10^{-31} \text{ kg}$

$$\lambda = h / m v$$

$$= (6.63 \times 10^{-34})$$

$$/ (9.11 \times 10^{-31} \times 5.4 \times 10^6)$$

$$\text{approx } 1.35 \times 10^{-10} \text{ m} = 1.35 \text{ \AA}$$

(b) $m v = 0.15 \times 20 = 3.0 \text{ kg m/s}$

$$\lambda = 6.63 \times 10^{-34} / 3.0$$

$$\text{approx } 2.2 \times 10^{-34} \text{ m (TINY!)}$$

(c) $m_p = 1.67 \times 10^{-27} \text{ kg}$

$$\lambda = h / m_p v$$

$$= 6.63 \times 10^{-34} / (1.67 \times 10^{-27} \times 10^5)$$

$$\text{approx } 3.97 \times 10^{-12} \text{ m} = 0.04 \text{ \AA}$$

Lesson : For the ball, $\lambda \ll$ atomic size \rightarrow wave nature undetectable.

For electrons, wave nature is dominant.

Davisson - Germer Experiment

1927 - first direct experimental proof of electron waves, by Davisson and Germer.

Idea

If electrons have wave nature, a beam should show diffraction maxima off a crystal lattice, like X-rays do.

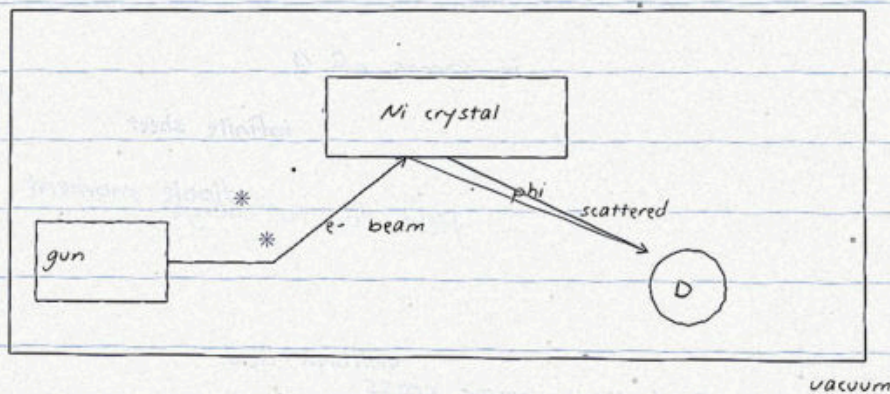


Fig. Davisson - Germer setup (top view)

Procedure

Vary V (accel.) and angle ϕ . • Measure intensity of scattered electrons at D.

Davisson - Germer Result

Sharp peak in scattered intensity at :

$$V = 54 \text{ V} , \quad \phi = 50 \text{ degrees.}$$

de Broglie prediction

$$\text{At } V = 54 \text{ V} ,$$

$$\lambda = 12.27 / \sqrt{54} \text{ \AA}$$

$$\text{approx } 1.67 \text{ \AA.}$$

Bragg prediction

For Ni crystal , $d = 0.91 \text{ \AA}$, $\phi = 50$.

$$\lambda = 2 \cdot d \sin(\phi / 2)$$

$$= 2 (0.91) \sin 25$$

$$\text{approx } 1.65 \text{ \AA} \quad (\text{matches!})$$

$$\lambda \text{ (de Broglie)} = \lambda \text{ (Bragg)}$$

← wave

← nature !

Significance

- * Confirmed de Broglie hypothesis for e-.
- * Foundation of ~~matter~~ wave mechanics .
- * Led to development of the electron microscope (G. P. Thomson, 1928).
- * Nobel Prize to Davisson & Thomson (1937).

Heisenberg's Uncertainty

Consequence of wave nature : one cannot fix both the position AND the momentum of a particle with arbitrary precision.

$$\Delta x \cdot \Delta p \geq h / (4 \pi)$$

*

(Often written $\Delta x \cdot \Delta p \geq \hbar / 2$)

Physical picture

A perfectly localized particle ($\Delta x \rightarrow 0$) needs a wave packet of infinitely many wavelengths \rightarrow momentum totally uncertain.

A pure plane wave (one λ) has fixed p but is spread over all space ($\Delta x \rightarrow \infty$).

Numerical feel

Electron in an atom : $\Delta x \approx 10^{-10} \text{ m}$.

$$\Delta p \geq h / (4 \pi \Delta x)$$

$$\text{approx } 5 \times 10^{-25} \text{ kg m/s}$$

$$\Delta v \geq \Delta p / m \text{ approx } 6 \times 10^5 \text{ m/s.}$$

Huge - classical orbits are not meaningful.

Photon Energy - Practical Cases

Compute $E = h c / \lambda$ for typical light.

($h c$ approx 1240 eV nm - handy shortcut)

$$E \text{ (eV)} = 1240 / \lambda \text{ (in nm)} \quad \leftarrow \text{useful ?}$$

Worked numbers

Red light : $\lambda = 700 \text{ nm}$

$$E = 1240 / 700 \text{ approx } 1.77 \text{ eV}$$

Green : $\lambda = 550 \text{ nm}$

$$E = 1240 / 550 \text{ approx } 2.25 \text{ eV}$$

Violet : $\lambda = 400 \text{ nm}$

$$E = 1240 / 400 \text{ approx } 3.10 \text{ eV}$$

X-ray : $\lambda = 0.1 \text{ nm}$

$$E = 1240 / 0.1 = 12400 \text{ eV} = 12.4 \text{ keV}$$

Notice : visible light just clears W for Cs, K, Na ; but cannot eject from Cu, Pt.

X-rays easily eject - even from heavy metals with KE in keV range.

Photo - cell Applications

Device based on photoelectric effect :

- * Evacuated glass tube .
- * Concave Cs - coated cathode catches light.
- * Anode collects photoelectrons.
- * Series with battery + microammeter.

Uses

- * Automatic doors , burglar alarms.
- * Sound track readers in old cinema film.
- * Light meters in cameras.
- * Counters in solar / lunar probes.
- * Photo - diodes / photo - transistors.

Electron microscope

Uses matter waves of accelerated electrons.

Since $\lambda_e \ll \lambda_{\text{visible}}$, resolution is FAR better than an optical microscope.

At $V = 60 \text{ kV}$, λ approx 0.05 \AA .

Resolution few angstroms \rightarrow individual atoms can be imaged.

Also : electron diffraction for crystal study.

Chapter Summary - Formulas

Photoelectric effect

$$E = h \nu = h c / \lambda \quad (\text{photon energy})$$

$$p = h / \lambda \quad (\text{photon momentum})$$

$$h \nu = W + KE_{\max} \quad (\text{Einstein})$$

$$e V_0 = KE_{\max} \quad (\text{stopping pot.})$$

$$W = h \nu_0 \quad (\text{threshold})$$

$$V_0 = (h/e) \nu - W/e \quad (\text{linear graph})$$

Matter waves

$$\lambda = h / p = h / m v$$

$$\lambda = h / \sqrt{2 m KE}$$

$$\lambda = h / \sqrt{2 m q V}$$

$$\lambda (\text{\AA}) = 12.27 / \sqrt{V} \quad \text{for } e^-$$

Constants to remember

$$h = 6.626 \times 10^{-34} \text{ J s}$$

$$e = 1.602 \times 10^{-19} \text{ C}$$

$$m_e = 9.11 \times 10^{-31} \text{ kg}$$

$$c = 3 \times 10^8 \text{ m / s}$$

$$h c = 1240 \text{ eV nm} \quad (\text{shortcut})$$

$$1 \text{ eV} = 1.6 \times 10^{-19} \text{ J}$$

$$\text{Uncertainty : } \Delta x \cdot \Delta p \geq h / (4 \pi)$$

End of Chapter 11 - Dual Nature.